

# 7<sup>th</sup> International Workshop on Charm Physics

## Wayne State University, 20/5/2015



**Measurements of T-odd Observables**  
**Maurizio Martinelli (EPFL)**  
**on behalf the LHCb collaboration**

# Outline

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## Theoretical introduction

- $CPV$  and  $T$ -odd correlations

## Previous searches

- $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$  and  $D_{(s)}^+ \rightarrow K^+ K_s^0 \pi^+ \pi^-$  decays at FOCUS, BaBar

## Search at LHCb

- $CPV$  search using  $T$ -odd correlations in  $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$  decays

## Conclusions

# CP Violation in Charm Mesons Decays

## Small CKM contribution

G. Isidori et al., Phys. Lett. B711 (2012) 46

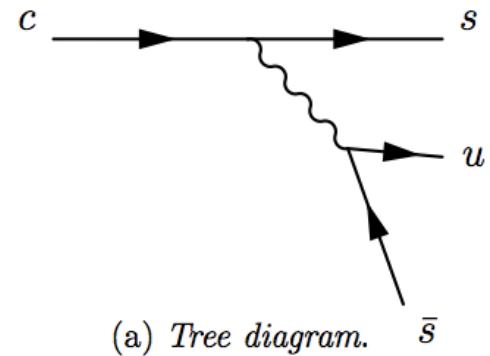
- CPV is expected to be small in the charm sector since tree-level topology dominates the decay at scales  $m_c < \mu < m_b$

## New Physics

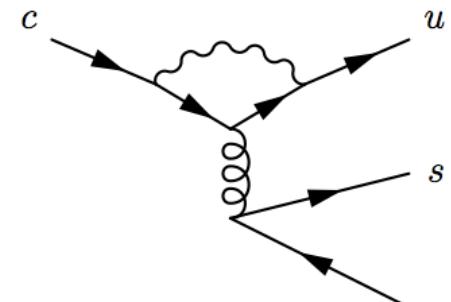
- May enhance this amplitude through the introduction of new processes and particles

## Recently

- Mixing in charm is established within SM expectation
- Experiments have recorded enough statistics to probe CPV with sensitivities approaching  $10^{-3}$
- If any NP effect is out there we should start to be able to see it



(a) Tree diagram.



(b) Penguin diagram.

# CP Violation Observables

## CPV in decay

$$\mathcal{A}_f \equiv \frac{\Gamma(M \rightarrow f) - \Gamma(\bar{M} \rightarrow \bar{f})}{\Gamma(M \rightarrow f) + \Gamma(\bar{M} \rightarrow \bar{f})}$$

- CPV is measured as the asymmetry between the decay rate of a meson and charge-conjugate state

$$\mathcal{A}_{f^\pm} = -\frac{2|a_1 a_2| \sin(\delta_2 - \delta_1) \sin(\phi_2 - \phi_1)}{|a_1|^2 + |a_2|^2 + 2|a_1 a_2| \cos(\delta_2 - \delta_1) \cos(\phi_2 - \phi_1)}$$

## CPV in mixing

$$\frac{q}{p} \neq 1, \quad \mathcal{A}_{SL}(t) \equiv \frac{d\Gamma/dt(\bar{M}_{\text{phys}}^0(t) \rightarrow l^+ X) - d\Gamma/dt(M_{\text{phys}}^0(t) \rightarrow l^- X)}{d\Gamma/dt(\bar{M}_{\text{phys}}^0(t) \rightarrow l^+ X) + d\Gamma/dt(M_{\text{phys}}^0(t) \rightarrow l^- X)}$$

- CPV measured from the mixing parameters

$$\mathcal{A}_{SL} = -\left| \frac{\Gamma_{12}}{M_{12}} \right| \sin(\phi_M - \phi_\Gamma)$$

## CPV in interference between decay and mixing

$$Im(\lambda_f) \neq 0, \quad \lambda_f \equiv \frac{q}{p} \frac{\bar{A}_f}{A_f}, \quad \mathcal{A}_{f_{CP}}(t) \equiv \frac{d\Gamma/dt(\bar{M}_{\text{phys}}^0(t) \rightarrow f_{CP}) - d\Gamma/dt(M_{\text{phys}}^0(t) \rightarrow f_{CP})}{d\Gamma/dt(\bar{M}_{\text{phys}}^0(t) \rightarrow f_{CP}) + d\Gamma/dt(M_{\text{phys}}^0(t) \rightarrow f_{CP})}$$

- CPV asymmetry is modified by mixing effects

$$\mathcal{A}_{f_{CP}}(t) = \eta_f \sin(\phi_M + 2\phi_f) \sin(\Delta m t)$$

$\phi$ : weak phases, from CKM  
 $\delta$ : unitarity (strong) phases

## A different point of view

- Describe invariant matrix element in the most general way (quasi two-body)  $B(p) \rightarrow V_1(k, \epsilon_1)V_2(q, \epsilon_2)$

$$M = a\epsilon_1 \cdot \epsilon_2 + \frac{b}{m_1 m_2} (p \cdot \epsilon_1)(p \cdot \epsilon_2) + i \frac{c}{m_1 m_2} \epsilon^{\alpha\beta\mu\nu} \epsilon_{1\alpha} \epsilon_{2\beta} k_\mu p_\nu$$

**S+D+P**

$$\begin{aligned} a &= \sum_j |a_j| e^{i(\delta_{sj} + \phi_{sj})}; & b &= \sum_j |b_j| e^{i(\delta_{dj} + \phi_{dj})}; & c &= \sum_j |c_j| e^{i(\delta_{pj} + \phi_{pj})} \\ \bar{a} &= \sum_j |a_j| e^{i(\delta_{sj} - \phi_{sj})}; & \bar{b} &= \sum_j |b_j| e^{i(\delta_{dj} - \phi_{dj})}; & \bar{c} &= \sum_j |c_j| e^{i(\delta_{pj} - \phi_{pj})} \end{aligned}$$

- A triple-product correlation arises in  $|M|^2$  from terms involving the c amplitude ( $\text{Im}[ac^*]$ ,  $\text{Im}[bc^*]$ ) wrt  $\vec{k} \cdot (\vec{\epsilon}_1 \times \vec{\epsilon}_2)$

$$\begin{aligned} A_B &= \frac{\Gamma(k \cdot \epsilon_1 \times \epsilon_2 > 0) - \Gamma(k \cdot \epsilon_1 \times \epsilon_2 < 0)}{N_B} \\ &\approx \Im(ac^*) = |ac| e^{i(\delta_s - \delta_p)} e^{i(\phi_s - \phi_p)} = |ac| \sin(\Delta\delta + \Delta\phi) \end{aligned}$$

- The same observable on the charge-conjugate decay gives

$$A_{\bar{B}} \approx |ac| e^{i(\delta_s - \delta_p)} e^{-i(\phi_s - \phi_p)} = |ac| \sin(\Delta\delta - \Delta\phi)$$

- That allows the definition of the CPV observable

$$a_{CP}^{T-\text{odd}} = A_B + A_{\bar{B}} \approx \cos \Delta\delta \sin \Delta\phi$$

## They are complementary

- The only difference is in the unitarity phases that enter differently in the game

$$a_{CP} \propto \sin \Delta\delta \sin \Delta\phi$$
$$a_{CP}^{T-\text{odd}} \propto \cos \Delta\delta \sin \Delta\phi \quad (*)$$

- $a_{CP}$  is more sensitive to CPV when the difference in the strong phases is large
- $a_{CP}^{T-\text{odd}}$  is more sensitive to CPV when the difference in the strong phases between the interfering amplitudes is small
- Datta and London demonstrated that a TP asymmetry can be also built with interference between decay and mixing, but it is proportional to  $\sin\Delta\delta$  as well.

(\*) **Caveat:** in  $a_{CP}$  the two phases are from different diagrams, in  $a_{CP}^{T\text{-odd}}$  from different spin contributions

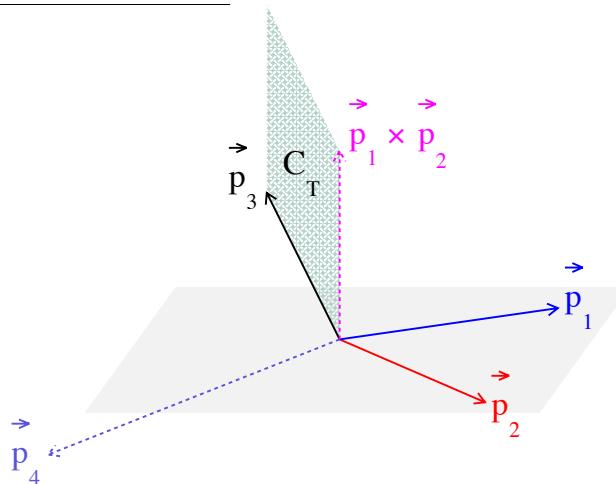
# Experimental Technique

## Defining a T-odd observable

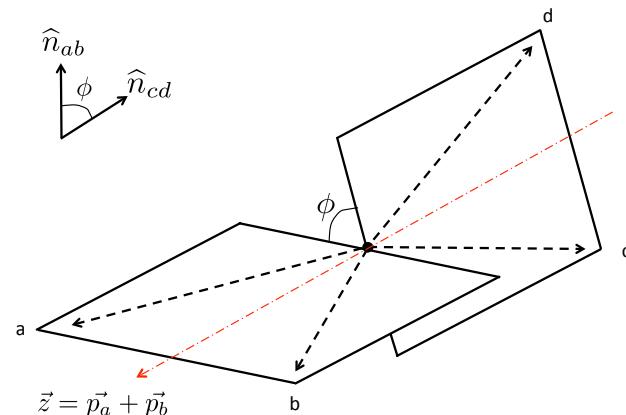
- One needs at least 3 independent momentum or spin variables

### 4-body decay

mother rest frame



$$C_T = (\vec{p}_1 \times \vec{p}_2) \cdot \vec{p}_3$$



$$\sin \phi = (\hat{n}_{ab} \times \hat{n}_{cd}) \cdot \hat{z}$$

- Momenta can be also used to define angles

# T-odd Correlation Asymmetry

## Asymmetries

- Two asymmetries are measured separately on the particle and charge-conjugate decays

$$A_T = \frac{\Gamma(C_T > 0) - \Gamma(C_T < 0)}{\Gamma}$$
$$\bar{A}_T = \frac{\bar{\Gamma}(-\bar{C}_T > 0) - \bar{\Gamma}(-\bar{C}_T < 0)}{\bar{\Gamma}}$$

- The  $CP$ -violating asymmetry is

$$a_{CP}^{T-\text{odd}} = \frac{1}{2}(A_T - \bar{A}_T)$$

# Results From Previous Experiments

## Used prompt $D^0$ and $D_{(s)}^+$ decays

- **FOCUS (2005)  $N_{\text{ev}} \sim 1\text{k}$**

Link et al., Phys. Lett. B662 (2005) 239

$$a_{CP}^{T-\text{odd}}(D^0) = (1.0 \pm 5.7(\text{stat}) \pm 3.7(\text{syst}))\%$$

$$a_{CP}^{T-\text{odd}}(D^+) = (2.3 \pm 6.2(\text{stat}) \pm 2.2(\text{syst}))\%$$

$$a_{CP}^{T-\text{odd}}(D_s^+) = (-3.6 \pm 6.7(\text{stat}) \pm 2.3(\text{syst}))\%$$

- **BaBar (2010-2011)  $N_{\text{ev}} \sim 50\text{k}$**

del Amo Sanchez et al., Phys. Rev. D81 (2010) 111103(R)  
Lees et al., Phys. Rev. D84 (2011) 031103(R)

$$a_{CP}^{T-\text{odd}}(D^0) = (1.0 \pm 5.1(\text{stat}) \pm 4.4(\text{syst})) \times 10^{-3}$$

$$a_{CP}^{T-\text{odd}}(D^+) = (-12.0 \pm 10.0(\text{stat}) \pm 4.6(\text{syst})) \times 10^{-3}$$

$$a_{CP}^{T-\text{odd}}(D_s^+) = (-13.6 \pm 7.7(\text{stat}) \pm 3.4(\text{syst})) \times 10^{-3}$$

- **BaBar provided significant statistical improvement (x10)**

## Semileptonic B decay

- $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$  from semileptonic B decays, tagged from muon charge  
 $B \rightarrow D^0 \mu^- X, D^0 \rightarrow K^+ K^- \pi^+ \pi^-$
- Clean sample

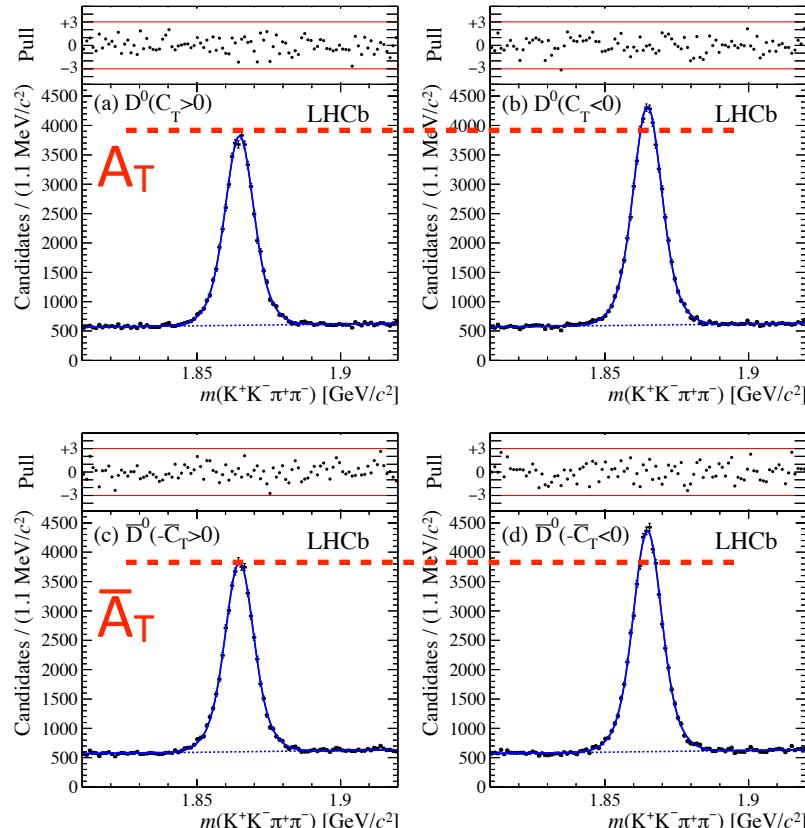
## Data Sample

- 2011+2012:  $3\text{fb}^{-1}$

## Fit Model

- Samples simultaneously fit to a model of two Gaussian distributions over an exponential shape
- Asymmetry parameters extracted from the fit

$B \rightarrow D^0 \mu^- X, D^0 \rightarrow K^+ K^- \pi^+ \pi^-$



$3\text{fb}^{-1}: N_{\text{ev}} \sim 170\text{k}$

## Three Measurements

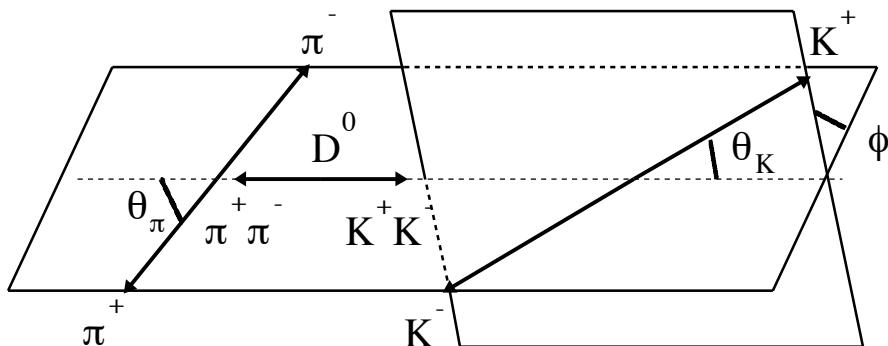
### 1. Integrated

$$a_{CP}^{T\text{-odd}}(D^0) = (1.8 \pm 2.9(\text{stat}) \pm 0.4(\text{syst})) \times 10^{-3}$$

### 2. Bins of phase-space

No significant deviation from 0 observed

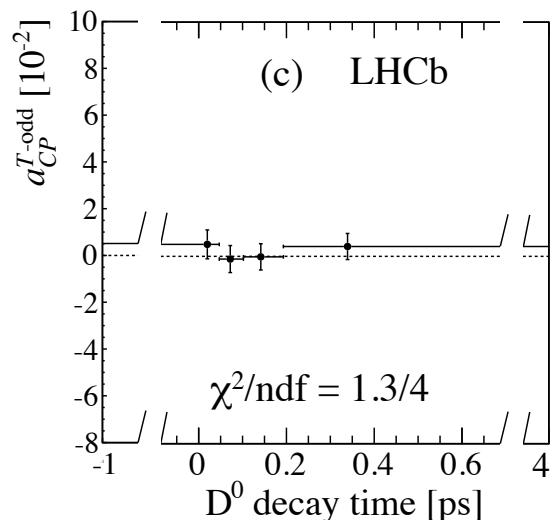
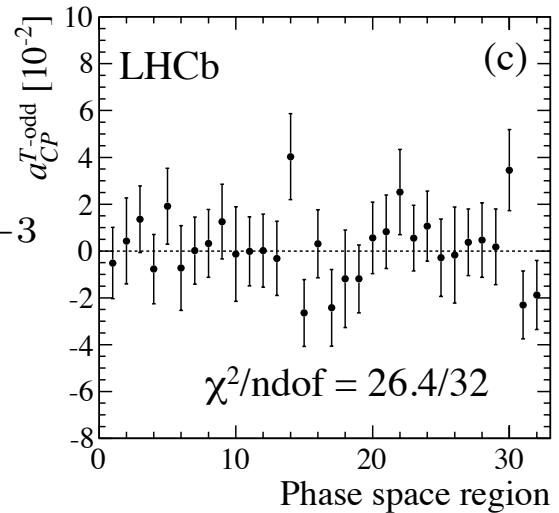
$CP$  conservation tested with  $P(\chi^2)=74\%$



### 3. Bins of $D^0$ decay time

No significant deviation from 0 observed

$CP$  conservation tested with  $P(\chi^2)=83\%$

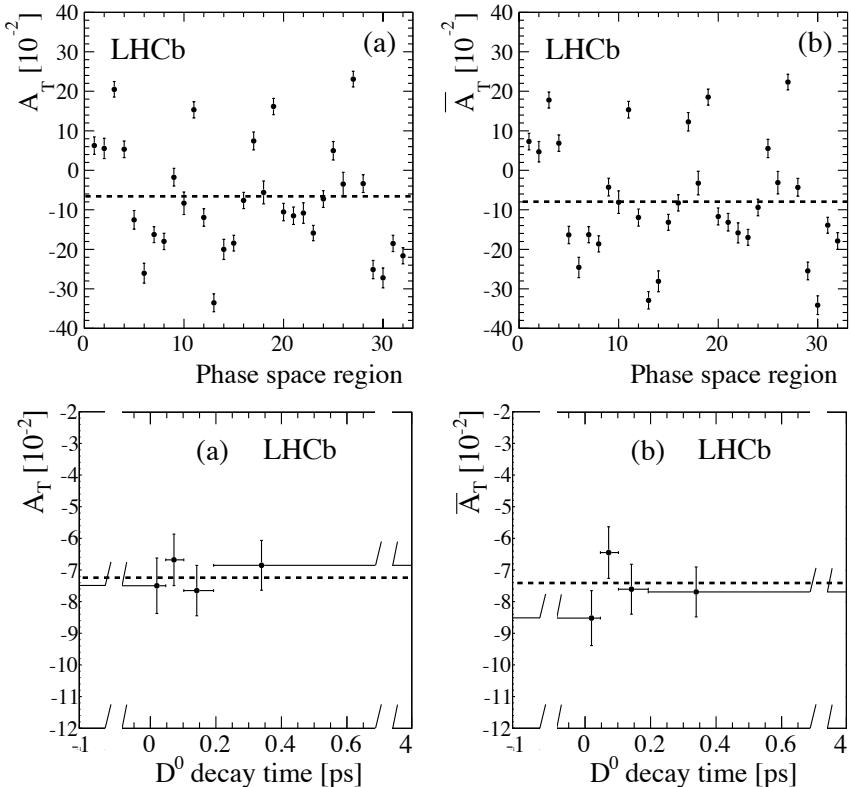


## FSI Effects?

- It's possible that FSI are producing effects in all the three measurements
- Significant differences in bins of phase space
- Average consistent wrt  $D^0$  decay time
- Wide spectrum of resonances and rescattering among the final state particles

$B \rightarrow D^0 \mu X, D^0 \rightarrow K^+ K^- \pi^+ \pi^-$

Local asymmetries up to 30%

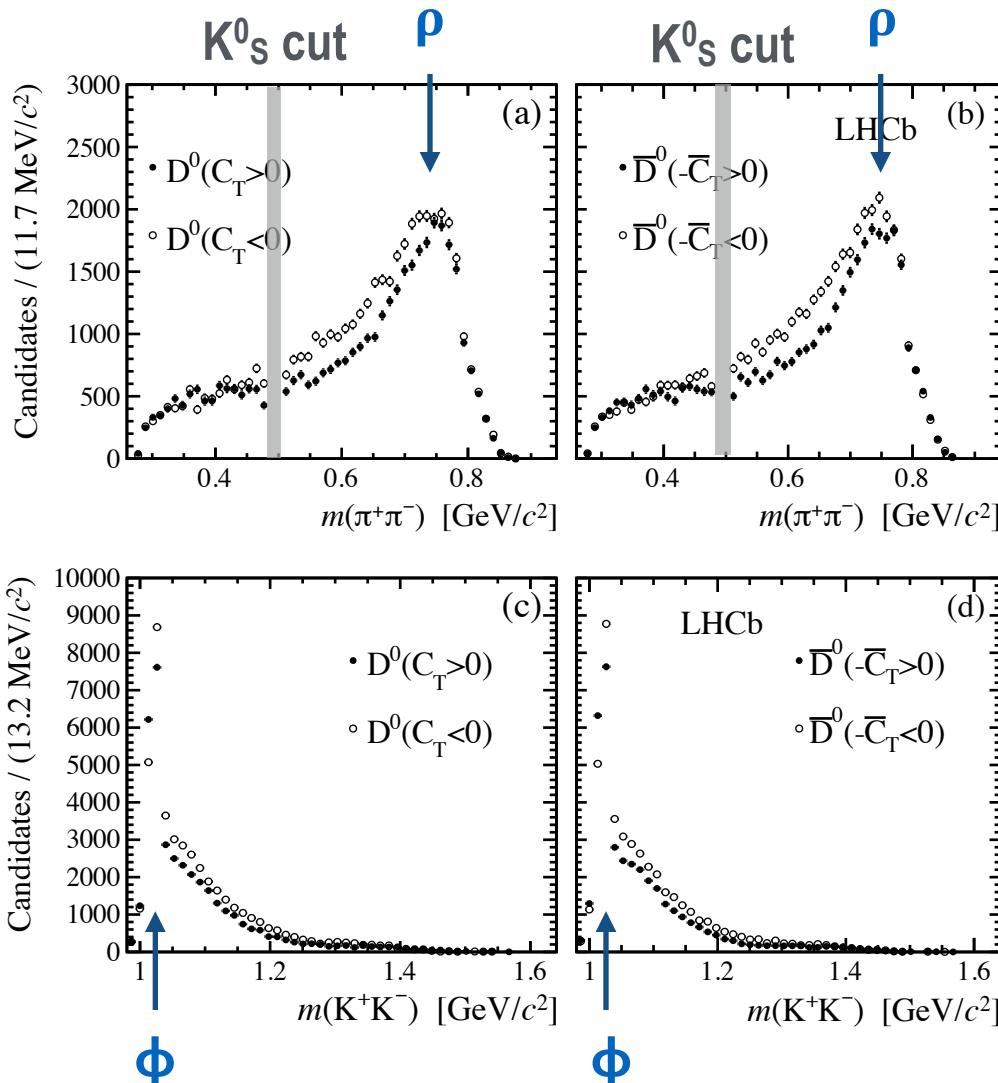


$$A_T(D^0) = (-71.8 \pm 4.1(\text{stat}) \pm 1.3(\text{syst})) \times 10^{-3}$$

$$\bar{A}_T(D^0) = (-75.5 \pm 4.1(\text{stat}) \pm 1.2(\text{syst})) \times 10^{-3}$$

## Resonant structure in $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$

- Clear evidence for  $\phi$  and  $\rho$  resonances
- Significant difference in the distributions vs  $C_T$
- Visible effects in angular variables as well
- $D^0 \rightarrow K^0_s K^+ K^-$  removed by  $\pi^+ \pi^-$  invariant mass cut



## Reconstruction Efficiency ☺

- Does not affect at all the result:  $A_T$  and  $\bar{A}_T$  asymmetries are calculated separately on the same final state

## Particle Identification ☺

- The same considerations apply to particle identification

## $C_T$ Resolution 🎉

- Estimated accurately from Monte Carlo, almost cancels in  $a_{CP}^{T\text{-odd}}$

## Peaking Backgrounds under $D^0/\bar{D}^0$ signal 🎉

- Any contamination affects the asymmetry as  $A \rightarrow A(1 - f) + f A^d$  ← very small effect  
 $f$  - contamination fraction;  $A^d$  - asymmetry of the contamination sample

## Flavour Mistag 🎉

- Considering the events with flavour mistag as a contamination  $a_{CP}^{T\text{-odd}} \rightarrow a_{CP}^{T\text{-odd}} - \Delta\omega/2(A_T + \bar{A}_T)$   
 $\Delta\omega = \omega^+ - \omega^-$  — difference among the mistag probabilities, measured from control samples  
 $B \rightarrow D^{*+} \mu^- X, (D^{*+} \rightarrow D^0 \pi^+, D^0 \rightarrow K^+ K^- \pi^+ \pi^-); B \rightarrow D^0 \mu^- X (D^0 \rightarrow K^- \pi^+ \pi^+ \pi^-)$

## Detector bias 🎉

- Conservative estimate from control sample of CF  $D^0 \rightarrow K^- \pi^+ \pi^+ \pi^-$

## Systematic uncertainty estimates

Contribution	$\Delta A_T(\%)$	$\Delta \bar{A}_T(\%)$	$\Delta a_{CP}^{T\text{-odd}}(\%)$
Prompt background	±0.09	±0.08	±0.00
Detector bias	±0.04	±0.04	±0.04
$C_T$ resolution	±0.02	±0.03	±0.01
Fit Model	±0.01	±0.01	±0.01
Flavor misidentification	±0.08	±0.07	±0.00
Total	±0.13	±0.12	±0.04

## Results

- CPV is searched for in  $D^0 \rightarrow K^+ K^- \pi^+ \pi^-$  using:
  1. A measurement integrated over the phase space
  2. Measurements in different regions of the phase space
  3. Measurements as a function of the  $D^0$  decay time
- No CPV found
- Interesting information in the phase space of the decay: local  $A_T$  asymmetries up to 30%  
These results are interpreted as possible effects of FSI produced by the rich resonant structure of the decay

1<sup>st</sup> time!

## Remarks

- Systematic uncertainties are found to be very small (as expected) in these observables  
High statistics control samples, toy studies

## LHCb Potential

- The full LHCb potential is not exploited yet  
A similar measurement on  $D^{*+} \rightarrow D^0 (\rightarrow K K \pi \pi) \pi^+$  decays is ongoing and will roughly double the statistics
- Run2 just behind the corner  
Improved trigger and selections promise 2.5 data per  $\text{fb}^{-1}$  wrt to Run1
- The rich resonant structure awaits detailed investigation  
Model-dependent and -independent approaches under study
- Other triple product correlations can be studied  
See talk tomorrow from A. Bevan ([hep-ph/1408.3813](#))

# Conclusions

## Alternative and Complementary CPV Tests

- Searches for CPV through asymmetries in  $T$ -odd moments are alternative and complementary to “standard” CPV measurements
- Applicable to many possible particle decays

## Low Systematics

- Analyses have demonstrated that the systematic uncertainties are very small

## Outlook

- Given the very low systematic uncertainties, such measurements could become extremely competitive at LHCb ( $10\text{fb}^{-1}$ ) and at future experiments (Belle-II, LHCb Upgrade,...)